

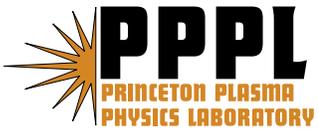
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# Observation of magnetocoriolis waves in a liquid metal Taylor-Couette experiment

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## Abstract

The first observation of fast and slow magnetocoriolis (MC) waves in a laboratory experiment is reported. Rotating nonaxisymmetric modes arising from a magnetized turbulent Taylor-Couette flow of liquid metal are identified as the fast and slow MC waves by the dependence of the rotation frequency on the applied field strength. The observed slow MC wave is marginally damped but will become destabilized by the magnetorotational instability with a modest increase in rotation rate.

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Hydromagnetic waves are a ubiquitous feature of both geophysical and solar dynamo models as well as models of astrophysical accretion disks. Such waves were first observed in experiments<sup>1</sup> with liquid metals using a sufficiently strong magnetic field that the Lorentz force could act like tension on a string and support Alfvén waves<sup>2</sup>. When the liquid is in rapid rotation, these waves are modified by the Coriolis force. The resultant magnetocoriolis waves<sup>3</sup> are a hybrid of Alfvén waves and inertial waves<sup>4</sup>. Magnetocoriolis waves (MC waves) are used to explain the secular variation of the Earth’s magnetic field over the course of hundreds of years<sup>5</sup> and the redistribution of angular momentum in the Sun<sup>6</sup>. They are a special case of the more general Magnetic Archimedes Coriolis (MAC) waves from dynamo theory<sup>7</sup>.

Despite the importance of MC waves in rotating conducting fluids and plasmas there is scant experimental evidence of their existence and of their relationship to various important astrophysical phenomena such as the dynamo or the Magnetorotational Instability (MRI)<sup>8</sup>. Recent experiments<sup>9</sup> on the DTS liquid sodium spherical Couette device and simulations<sup>10</sup> have found several different types of hydromagnetic waves, but there is ambiguity about their identification as MC waves. A similar experiment in Maryland found inertial waves<sup>11</sup>, but the applied fields were too weak to observe Lorentz force effects.

In this Letter, we report the first clear identification of the combined fast and slow MC waves in a laboratory experiment. Through measurements of the radial magnetic field in a liquid metal Taylor-Couette flow, we observe two rotating modes that follow the rotation speeds expected for the fast and slow MC-wave. We also demonstrate through a local stability analysis that with the addition of sufficient flow shear, the slow MC-wave can be destabilized to produce the MRI. Using the observed frequencies of the waves we infer from the local dispersion relation that the modes are damped and obtain a method of determining the threshold for the MRI.

The Princeton MRI experiment is designed to study the stability of a rotating sheared flow of liquid metal with an applied magnetic field coaxial with the rotation axis. The apparatus has been described elsewhere<sup>12</sup> and has already demonstrated the ability to generate high Reynolds number shear flow in water with angular momentum flux comparable to viscous transport (a null result in trying to demonstrate subcritical hydrodynamic instability<sup>13</sup>). The volume between the concentric rotating cylinders is filled with GaInSn, a gallium eutectic alloy<sup>14</sup>. The dimensions of the experiment and properties of the liquid metal are given in Tab. I.

The equations describing the evolution of a rotating shear flow with a background magnetic

Experimental Parameters	Symbol	Units	Value
Height	$h$	cm	27.9
Inner Cylinder radius	$r_1$	cm	7.06
Outer Cylinder radius	$r_2$	cm	20.30
Density	$\rho$	g/cm <sup>3</sup>	6.36
Kinematic viscosity	$\nu$	cm <sup>2</sup> /s	$2.98 \times 10^{-3}$
Magnetic diffusivity	$\eta$	cm <sup>2</sup> /s	$2.57 \times 10^3$

TABLE I: Physical parameters of the apparatus<sup>12</sup> and liquid metal<sup>14</sup>.

field are given by the magnetohydrodynamic (MHD) equations in a rotating frame:

$$\begin{aligned} \frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} + 2\mathbf{\Omega} \times \mathbf{v} &= -\nabla P + \frac{1}{\mu_0 \rho} (\mathbf{B} \cdot \nabla) \mathbf{B} + \nu \nabla^2 \mathbf{v} \\ \frac{\partial \mathbf{B}}{\partial t} &= \nabla \times (\mathbf{v} \times \mathbf{B}) + \eta \nabla^2 \mathbf{B} \\ P &= \frac{p}{\rho} + \frac{1}{2} \frac{B^2}{\mu_0 \rho} - \frac{1}{2} |\mathbf{\Omega} \times \mathbf{r}|^2 \end{aligned}$$

where  $\mathbf{v}$  and  $\mathbf{B}$  are solenoidal fields representing the velocity and magnetic field,  $\mathbf{\Omega}$  is the angular velocity, and  $p$  is the pressure. The generalized pressure  $P$  incorporates the magnetic and centrifugal pressure terms.

The background field and angular velocity are given by  $\mathbf{B}_0 = B_0 \hat{\mathbf{z}}$  and  $\mathbf{\Omega} = \Omega \hat{\mathbf{z}}$ . If we assume harmonic perturbations of the velocity and magnetic field and linearize the resulting equations, we obtain the dispersion relation in cylindrical coordinates<sup>15</sup>:

$$\begin{aligned} &(\bar{\omega} - i\gamma_\eta)^2 [(\bar{\omega} - i\gamma_\nu)(\bar{\omega} - i\gamma_\eta) + \omega_A^2]^2 \\ &+ 2\zeta\Omega^2 (\bar{\omega} - i\gamma_\eta)^4 (k_z/k)^2 - 2(2 - \zeta)\Omega^2 \omega_A^2 (\bar{\omega} - i\gamma_\eta) (k_z/k)^2 \\ &- \omega_R [(\bar{\omega} - i\gamma_\nu)(\bar{\omega} - i\gamma_\eta) + \omega_A^2] [(\bar{\omega} - i\gamma_\eta)^2 + \omega_A^2] = 0 \end{aligned} \quad (1)$$

where  $\bar{\omega} = \omega - m\Omega$  is the Doppler-shifted complex frequency,  $\gamma_\eta = \eta k^2$  and  $\gamma_\nu = \nu k^2$  are the resistive and viscous damping rates,  $\omega_A = k_z B_0 / \sqrt{\mu_0 \rho}$  is the Alfvén frequency, and  $\omega_R = (\zeta - 2)\Omega k_r k_\theta / k^2$  is the Rossby wave frequency<sup>16</sup>. We have quantified the flow shear by introducing the vorticity parameter

$$\zeta(r) = |\nabla \times r\Omega(r)\hat{\phi} / \Omega(r)| = (1/r\Omega)\partial(r^2\Omega)/\partial r. \quad (2)$$

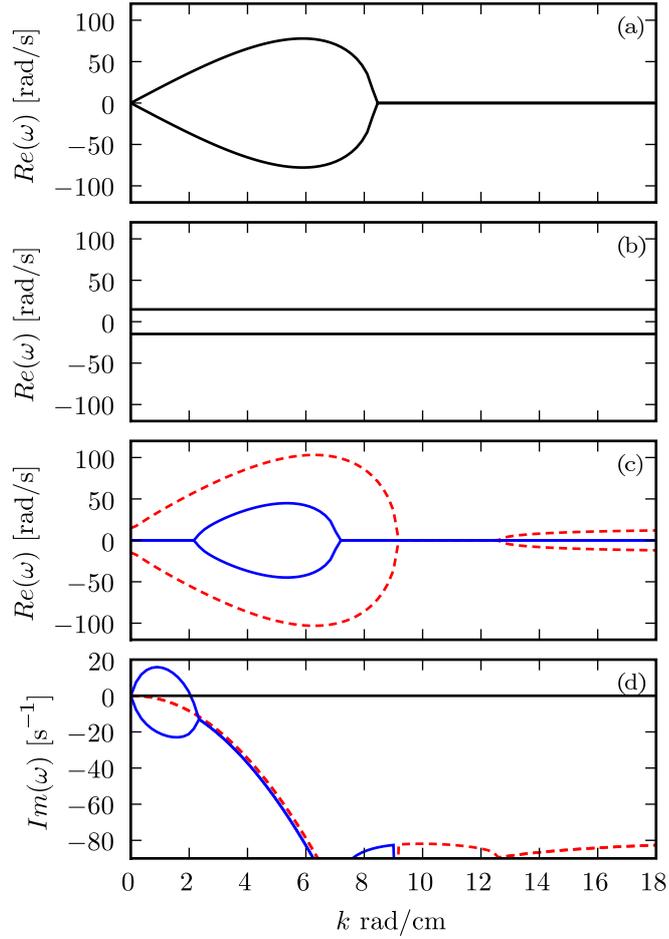


FIG. 1: The real part of the dispersion relation for (a) damped Alfvén waves as given in Eq. 3, (b) inertial waves as given in Eq 4, and (c) the fast (dashed red line) and slow MC waves (solid blue line). (d) The growth rate of the fast and slow MC waves. Due to resistive dissipation, the MC waves collapse to the inertial wave dispersion at high  $k$ . When there is sufficient flow shear, the slow MC-wave becomes unstable at low  $k$ . The parameters for the dispersion relation are  $B_0 = 4$  kG,  $\Omega = 42$  rad/s, and  $\zeta = 0.25$  and the fluid properties are given in Tab. I.

Note that  $\zeta$  is only constant if the rotation profile follows a power-law dependence  $\Omega(r) \propto r^{\zeta-2}$ . For uniform rotation with no shear  $\zeta = 2$ . The Rayleigh criterion<sup>17</sup>, which governs hydrodynamic centrifugal stability of rotating shear flow, is given by  $\zeta \geq 0$ .

We can gain insight into the basic waves for this dispersion relation by examining limiting

cases. In the absence of rotation, the dispersion relation reduces to

$$(\omega - i\gamma_v)(\omega - i\gamma_\eta) - \omega_A^2 = 0 \quad (3)$$

which describes the damped shear Alfvén wave. These waves have a transverse polarization due to the incompressibility of the fluid (in a compressible fluid there is also a longitudinal magnetosonic wave). The complex frequency is  $\omega = i(\gamma_v + \gamma_\eta)/2 \pm \sqrt{\omega_A^2 - (\gamma_v - \gamma_\eta)^2}$  which shows that the Alfvén wave is viscously and resistively damped and has a real frequency of  $\pm\omega_A$  in the absence of dissipation. Variations in the flow along the direction of the magnetic field tend to be eliminated in a highly resistive fluid (such as a liquid metal) due to this damping.

Assuming rotation without shear and without an applied magnetic field, the dispersion relation reduces to

$$(\omega - i\gamma_v)^2 + (2\Omega k_z/k)^2 = 0 \quad (4)$$

which describes inertial waves. Inertial waves also have a transverse polarization but are peculiar in that the restoring force, provided by the Coriolis effect, acts orthogonally to the displacement. The resulting motion of a displaced fluid element is circular precession. The complex frequency is  $\omega = i\gamma_v \pm 2\Omega k_z/k$  so the wave is viscously damped with a real frequency between  $\pm 2\Omega$ . Note that there is no dependence on the wavenumber. Akin to Alfvén waves, inertial waves homogenize the flow along the axis of rotation due to viscous damping, consistent with the Taylor-Proudman Theorem<sup>4</sup>.

Together these waves make up the hybrid magnetocoriolis wave<sup>3</sup>. Since there are two restoring forces acting on displaced fluid elements, there are two possible situations. The Lorentz and Coriolis forces may act together, stiffening the system and producing the higher frequency fast wave, or the two forces may oppose one another to produce the lower frequency slow wave. The requirement for observing strong rotational effects on the Alfvén wave is  $r\Omega\sqrt{\mu_0\rho}/B_0 \sim 1$ . The Alfvén frequency experiences a splitting due to the breaking of the degeneracy of the two roots by the presence of rotation. When the flow has sufficient shear, the slow MC-wave becomes stationary and unstable at low  $k$  as seen in Fig. 1. This instability is the MRI which in addition to causing turbulent transport of angular momentum in accretion disks may also be related to geomagnetic jerks<sup>18</sup>.

The results reported here were obtained when an axial magnetic field was applied to turbulent rotating shear flow. The outer cylinder was kept stationary while the inner cylinder rotated. The rings at the end caps were coupled to either cylinder in what is referred to as the “split” configu-

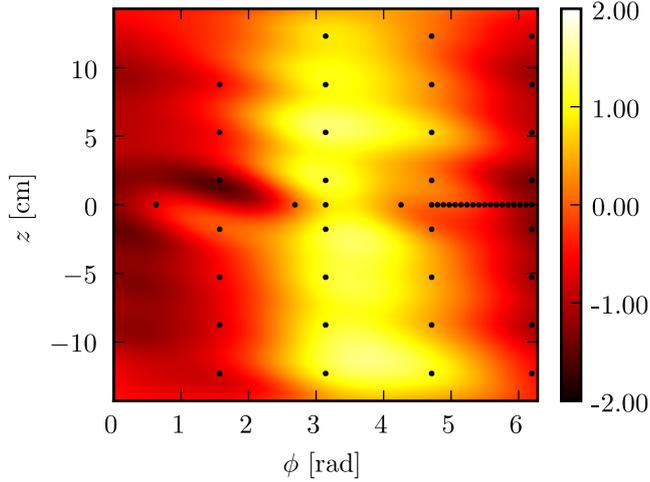


FIG. 2: A contour plot of the radial magnetic field near the surface of the outer cylinder constructed from an interpolation of data from the Mirnov coil array. The locations of the coils are depicted by dots and the contour levels are given in Gauss. The applied field was 4.3 kG.

ration<sup>12</sup>. The flow was hydrodynamically unstable since  $\zeta \leq 0$ . The inner cylinder and inner ring were set in rotation at 6.7 Hz and an axial magnetic field between 1.7 to 4.3 kG was applied to the turbulent flow. The induced radial magnetic field fluctuations were measured by an array of Mirnov coils positioned just outside the outer cylinder. A nonaxisymmetric mode was apparent from the coil array as seen in the snapshot of the field shown in Fig. 2. The signals from the array were then interpolated onto a regular grid and a 2D Fourier transform was applied to obtain the amplitude and phase of the two dominant modes shown in Fig. 3. Similar results are obtained when a more formal least-squares fitting technique is used<sup>26</sup>.

The mode rotation rates are measured by calculating the linear slope of the Fourier phase with time for each field strength. The results are shown in Fig. 4. The modes clearly rotate at different speeds and their rotation rates increase with magnetic field strength. By comparing the rotation rates with those of higher azimuthal harmonics obtained from the high density mid-plane coil array (seen as the series of dots at  $z = 0$  in Fig. 2) we find that the harmonics are not phase locked. Hence, the signal is not likely due to a passing vortex as was observed in Sisan et al.<sup>19</sup>.

A least squares fit of the observed mode rotation rates to the real frequencies of the fast and slow MC waves gives an estimate of the local wave vector components and the fluid rotation rate for a given shear. The vorticity parameter  $\zeta$  has no effect on the fit over the range where there is data and is assumed to be zero for marginal stability consistent with hydrodynamic observations<sup>20</sup>. The

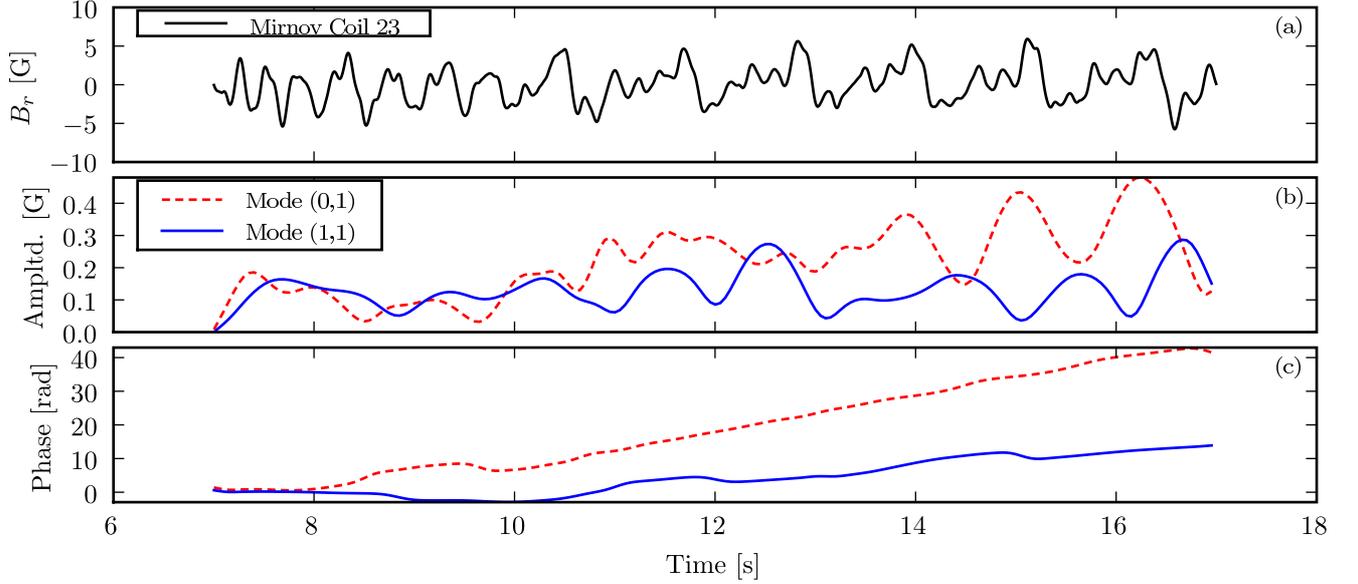


FIG. 3: (a) Magnetic field measured by a single Mirnov coil. (b) The time series of (0,1) and (1,1) mode amplitudes from a 2D Fourier decomposition of the radial magnetic field where the notation  $(n, m)$  describes the vertical mode number  $n$  and the azimuthal mode number  $m$ . (c) The time series of mode phases.

fit parameters result in a mode with  $k_z = \pi/2h$  and  $k_r = \pi/(r_2 - r_1)$ . The corresponding growth rate of the mode as a function of the applied magnetic field is shown in Fig. 4b. Aside from an expected Doppler shift for the nonaxisymmetric modes, the MC-wave model provides an excellent fit to the observations. From the growth rate, we find that the slow wave is marginally stable for 1 kG but is otherwise damped. Note as well that the slow wave is damped wherever it has nonzero frequency, a generic characteristic of the MRI. From these properties of the dispersion relation we can predict the necessary rotation rate and magnetic field strength to observe the MRI based on empirical observations of the damped waves.

Similar nonaxisymmetric waves have been observed in the PROMISE magnetized Taylor-Couette experiment<sup>21</sup> which may be related to our observed MC waves. Sisan et al.<sup>22</sup> also observed rotating nonaxisymmetric spherical harmonic patterns in a turbulent liquid sodium spherical Couette flow. Although they attribute their observation of rotating nonaxisymmetric modes to a bulk flow instability, the possibility that their observations are of a driven, damped mode should be examined.

Further work is required to identify the source of these damped waves. One conjecture is that the turbulent flow provides perturbations of a broad range of wavelengths, but that the geometry

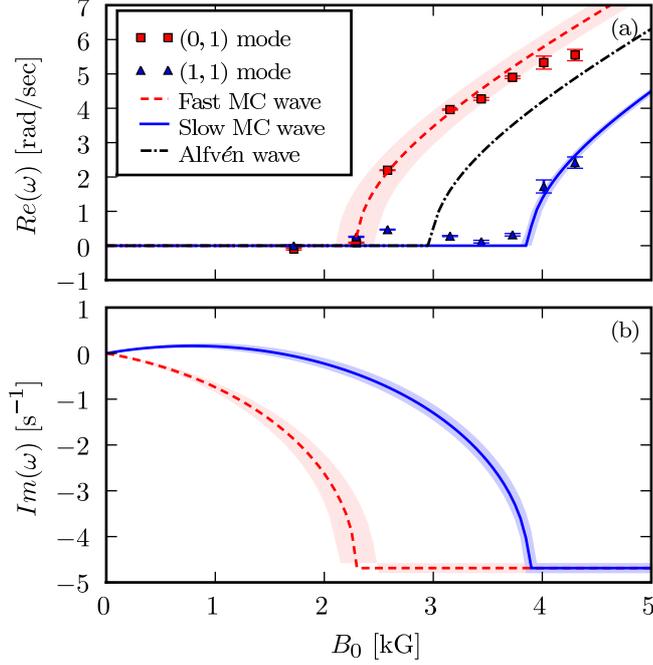


FIG. 4: The real frequency (a) and growth rate (b) determined by fitting the phase speed of the two non-axisymmetric modes observed for a range of applied magnetic field strengths to Eq. 1. Error bars reflect the uncertainty in the linear fit to the phase as a function of time. The dashed red and solid blue lines show the fit of the fast and slow MC waves from Eq. 1 with values of  $\zeta = 0$ ,  $k = 0.246 \pm 0.001 \text{ cm}^{-1}$ ,  $\theta = 1.336 \pm 0.005 \text{ rad}$  where  $k_z = k \cos \theta$ , and  $\Omega = 5.4 \pm 0.7 \text{ rad/s}$ . The shaded areas express the uncertainty in the fit.

of the vessel dictates which modes are realized<sup>23</sup>. Such is observed when a precessing top cap is used to drive inertial waves in a cylinder filled with water<sup>24</sup>. The observed waves in this case are cavity resonances driven by unstable flow at the boundary discontinuities and not an instability of the bulk flow. It is also possible that these damped waves are driven through nonlinear wave coupling<sup>25</sup>. If such is the case, then these damped waves may be a saturation mechanism for the MRI. It is due to the ambiguity of the source of these waves that we are continuing to pursue an observation of magnetically induced instability in a hydrodynamically quiescent flow.

In summary, we have observed rotating modes in a turbulent Taylor-Couette flow of liquid metal that we identify as the fast and slow MC-wave. We have elucidated the relationship between the slow MC-wave and the MRI and have deduced a method of determining the threshold for instability through observation of driven MC waves. MC waves will be important in identifying

the MRI in further experiments and may also play a role in saturation of MHD turbulence in rotating fluids.

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<sup>26</sup> These results will be published in a companion paper.



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